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**Progress towards the determination of an empirical flux equation for asymmetry-induced transport**<sup>1</sup> D. L. EGGLESTON, C. T. SMITH, Occidental College — In previous work<sup>2</sup> on asymmetry-induced transport, it was found useful to employ the hypothesis that the asymmetry frequency  $\omega$  and the plasma rotation frequency  $\omega_R$  always enter the physics in the combination  $\omega - l\omega_R$ , where  $l$  is the azimuthal mode number of the asymmetry. Flux data points satisfying the condition  $\omega - l\omega_R = 0$  were shown to satisfy the equation  $\Gamma_{sel} = -(B_0/B)^{1.33} D_0 [\nabla n_0 + f_0]$ , where  $B$  is the magnetic field,  $\nabla n_0$  is the radial density gradient, and  $B_0$ ,  $D_0$ , and  $f_0$  are empirical constants. The general flux equation was then constrained to be of the form  $\Gamma(\epsilon) = -(B_0/B)^{1.33} D(\epsilon) [\nabla n_0 + f(\epsilon)]$ , where  $\epsilon = \omega - l\omega_R$  and  $D(\epsilon)$  and  $f(\epsilon)$  are unknown functions. We now examine data points adjacent to the  $\epsilon = 0$  points and compare them to a first order expansion of  $\Gamma(\epsilon)$ . We find that a plot of  $d(\Gamma - \Gamma_{sel})/d\epsilon$  vs  $r$  changes sign at about the same radius as  $\nabla n_0 + f_0$ , and show that this implies that  $dD/d\epsilon(0) \neq 0$ . This, plus the requirement that  $D(\epsilon = 0) = D_0$ , restricts the form of  $D(\epsilon)$ . In particular, it excludes a dependence on  $\epsilon$  of the form found in resonant particle transport theory<sup>3</sup>, i.e.,  $D(\epsilon) \propto \exp(-C\epsilon^2)$ , with  $C$  a parameter.

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<sup>2</sup>D. L. Eggleston and J. M. Williams, Phys. Plasmas 15, 032305 (2008).

<sup>3</sup>D. L. Eggleston and T. M. O'Neil, Phys. Plasmas 6, 2699 (1999).

☐ Prefer Oral Session  
☒ Prefer Poster Session

Dennis L. Eggleston  
dleggles@oxy.edu  
Occidental College

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